

# Unitary evolutions

(recap / high speed overview of quantum gates)

→ Single example of type of quantum evolution

$$i\hbar \frac{\partial |\psi\rangle}{\partial t} = H |\psi\rangle$$

Schrödinger's Eq.

In quantum information we often abstract away from the underlying Hamiltonian & think entirely in terms of unitary operations applied to quantum states.

$$|\psi\rangle = \underbrace{e^{-iHt/\hbar}}_U |\psi\rangle$$

a unitary matrix  
 $U U^\dagger = I$

2 level system  
- basic building block of quantum computation

∴ Reversible  $U^\dagger(U|\psi\rangle) = |\psi\rangle$

& length preserving

$$|0\rangle = U|\psi\rangle$$

$$\langle 0|0\rangle = \langle \psi|U^\dagger U|\psi\rangle = \langle \psi|\psi\rangle$$

## ① Single qubit gates

Basic gates:

ie. basic set of unitary operations to apply to a single qubit

classical

- Identity  $\begin{matrix} |0\rangle \rightarrow |0\rangle \\ |1\rangle \rightarrow |1\rangle \end{matrix} \}$  trivial  $I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$
- NOT  $\begin{matrix} |0\rangle \rightarrow |1\rangle \\ |1\rangle \rightarrow |0\rangle \end{matrix}$   $NOT = X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

(check:  $X|0\rangle = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix} = |1\rangle$

$X|1\rangle = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} = |0\rangle$ )

Not to be confused with the Hamiltonian

→ • Hadamard  $H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$

QUANTUM  
Generates  
non-classical  
states

$$H|0\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) = |+\rangle$$

$$H|1\rangle = |-\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$$



## • Evolution under Pauli Matrix

It's helpful to know how a unitary on a single qubit affects a state on a Bloch sphere.

$$\text{Let } \rho = \frac{1}{2}(\mathbb{I} + \underline{r} \cdot \underline{\sigma})$$

$$\rho_{\text{out}} = U \rho U^\dagger = \frac{1}{2}(\mathbb{I} + \underbrace{U(\underline{r} \cdot \underline{\sigma})U^\dagger}_{\equiv \underline{S} \cdot \underline{\sigma}})$$

Claim:  $\underline{S} = \mathbf{O} \underline{r}$  where  $\mathbf{O}$  is an orthogonal matrix

Orthogonal matrices induce (length preserving) rotations on real vectors.  $\left\{ \begin{array}{l} \text{i.e. } \mathbf{O} \mathbf{O}^T = \mathbf{O}^T \mathbf{O} = \mathbb{I} \\ \text{Real}(\mathbf{O}) = \mathbf{O} \end{array} \right.$

(i.e. they are the real analogue of unitaries)

Proof:

To show that  $\mathbf{O}$  is orthogonal we just need to show

$$\text{that } |\underline{r}| = |\underline{S}|$$

$$\text{To do so, we note that } \text{Tr}[(\underline{S} \cdot \underline{\sigma})^2] = \text{Tr}\left[\left(\sum_i s_i \sigma_i\right)^2\right]$$

$$\begin{aligned}
 &= \text{Tr}(\sum_{ij} s_i s_j \sigma_i \sigma_j) \\
 &= \sum_{ij} s_i s_j \delta_{ij} \\
 &= \sum_i s_i^2 = |\underline{s}|^2
 \end{aligned}$$

Now we note that

$$\begin{aligned}
 |\underline{s}|^2 &= \text{Tr}((\underline{s} \cdot \underline{\sigma})^2) = \text{Tr}(U (\underline{r} \cdot \underline{\sigma}) U^\dagger)^2 = \text{Tr}(U (\underline{r} \cdot \underline{\sigma})^2 U^\dagger) = \sum_{ij} \text{Tr}(r_i r_j \sigma_i \sigma_j) \\
 &= \sum_i r_i^2 = |\underline{r}|^2
 \end{aligned}$$

$\therefore |\underline{s}| = |\underline{r}| \Rightarrow \underline{O}$  is orthogonal

To get a better handle on the nature of this rotation, note that:

Any unitary on a qubit can be written as

$$U = e^{i\theta \underline{n} \cdot \underline{\sigma}}$$

(if this isn't immediately obvious note that any unitary can be written as

$U = e^{-iHt}$  and  $H$  can always be expanded in the Pauli basis)

$H$  is now a Hamiltonian (sorry!)

which is equivalent to

$$U = e^{i\theta \underline{n} \cdot \underline{\sigma}} = \cos(\theta) I - i \sin(\theta) \underline{n} \cdot \underline{\sigma}$$

(if you've not shown this before it's worth proving it to yourself)

What is the effect of evolving  $|0\rangle\langle 0|$  under  $e^{-i\theta\sigma_x}$ ?

$$e^{-i\theta\sigma_x} |0\rangle = \cos\theta |0\rangle - i\sin(\theta) \sigma_x |0\rangle \\ = \cos\theta |0\rangle - i\sin\theta |1\rangle$$

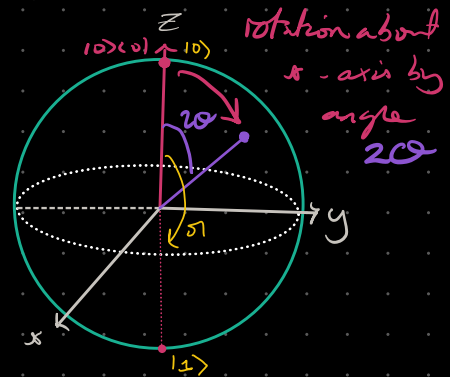
eg.  $e^{-i\frac{\pi}{2}\sigma_x} |0\rangle = \cos(\frac{\pi}{2}) |0\rangle - i\sin(\frac{\pi}{2}) |1\rangle$

Alt

$$= -i |1\rangle$$

unphysical global phase

Exercise: Show that  $e^{-i\theta \hat{n} \cdot \sigma}$  corresponds to a rotation about  $\hat{n}$  vector by an angle  $2\theta$ .



## 2 Two-Qubit Gates

Examples

- Can just apply multiple copies of a gate onto systems in parallel.

eg.  $X \otimes X |0\rangle \otimes |0\rangle = X |0\rangle \otimes X |0\rangle = |1\rangle \otimes |1\rangle$

Alternative common notation:  $\underbrace{X \otimes X}_{X^{\otimes 2}} \underbrace{|0\rangle \otimes |0\rangle}_{|0\rangle^{\otimes 2}}$

- More interesting: CNOT: "Apply NOT to 2nd Qubit iff 1st qubit is in state  $|1\rangle$ "

ie.  $\begin{aligned} \text{CNOT } |00\rangle &= |00\rangle \\ \text{CNOT } |01\rangle &= |01\rangle \\ \text{CNOT } |10\rangle &= |11\rangle \\ \text{CNOT } |11\rangle &= |10\rangle \end{aligned} \quad \equiv$

Truth table

Before	After
00	00
01	01
10	11

Explicit Matrix form:  $CNOT = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$

(Check this really does do)

$CNOT$  can generate entanglement.

eg. consider the sequence of gates

$$\begin{aligned} (CNOT)(H \otimes I) |00\rangle &= CNOT \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \otimes |0\rangle \\ &= CNOT \frac{1}{\sqrt{2}} (|00\rangle + |10\rangle) \\ &= \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle) \end{aligned}$$

non classical but unentangled state

Bell state

maximally entangled state

Bell states

$$\begin{aligned} |\phi^\pm\rangle &= \frac{1}{\sqrt{2}} (|00\rangle \pm |11\rangle) \\ |\psi^\pm\rangle &= \frac{1}{\sqrt{2}} (|01\rangle \pm |10\rangle) \end{aligned}$$

Qubits perfectly correlated with prob  $\frac{1}{2}$  both 00 &  $\frac{1}{2}$  11

Qubits perfectly anti-correlated with prob  $\frac{1}{2}$  first qubit 0 & second 1 & prob  $\frac{1}{2}$  first qubit 1 & second 0

How would you prepare the other Bell states using quantum gates?

### ③ Drawing circuits

A tidy way of visualising a series of gates

Best understood by example

