

# Unitary evolutions

(recap / high-speed overview of quantum gates)

→ Simple example of type of quantum evolution

$$i\hbar \frac{d|\psi\rangle}{dt} = H|\psi\rangle$$

Schrödinger's Eq.

$$|\psi\rangle = e^{-iHt}|\psi\rangle$$

a unitary matrix

$$U \quad UU^\dagger = I$$

In quantum information we often abstract away from the underlying Hamiltonian & think entirely in terms of unitary operations applied to quantum states

2 level system

- basic building

block of quantum computation

$$\therefore \text{Reversible} \quad U^\dagger(U|\psi\rangle) = |\psi\rangle$$

& length preserving

$$|\phi\rangle = U|\psi\rangle$$

$$\langle\phi|\phi\rangle = \langle\psi|U^\dagger U|\psi\rangle = \langle\psi|\psi\rangle$$

## ① Single qubit gates

Basic gates: i.e. basic set of unitary operations to apply to a single qubit

$$\left. \begin{array}{l} \bullet \text{ Identity} \quad |0\rangle \rightarrow |0\rangle \\ \quad \quad \quad |1\rangle \rightarrow |1\rangle \end{array} \right\} \text{trivial} \quad I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\left. \begin{array}{l} \bullet \text{ NOT} \quad |0\rangle \rightarrow |1\rangle \\ \quad \quad \quad |1\rangle \rightarrow |0\rangle \end{array} \right. \quad \text{NOT} = X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$(\text{check: } X|0\rangle = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix} = |1\rangle)$$

$$X|1\rangle = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} = |0\rangle)$$

Not to be confused with the Hamiltonian

$$\left. \begin{array}{l} \bullet \text{ Hadamard} \quad H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \end{array} \right.$$

Quantum  
Generated  
non-classical  
state

$$H|0\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) = |+\rangle$$

$$H|1\rangle = |-\rangle = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

### • Evolution under Pauli Matrix

It's helpful to know how a unitary on a single qubit effects a state on a Bloch sphere.

$$\text{Let } \rho = \frac{1}{2} (\mathbb{I} + \underline{\sigma} \cdot \underline{\sigma})$$

$$P_{\text{out}} = U \rho U^\dagger = \frac{1}{2} (\mathbb{I} + \underbrace{U (\underline{\sigma} \cdot \underline{\sigma}) U^\dagger}_{\equiv S \cdot \underline{\sigma}})$$

Claim:  $\underline{S} = \underline{\sigma} \cdot \underline{\sigma}$  where  $\underline{\sigma}$  is an orthogonal matrix

Orthogonal matrices induce (length preserving) rotations on real vectors.  $\left\{ \begin{array}{l} \text{i.e. } \underline{\sigma} \underline{\sigma}^T = \underline{\sigma}^T \underline{\sigma} = \mathbb{I} \\ \text{Real}(\underline{\sigma}) = \underline{\sigma} \end{array} \right.$

(i.e. they are the real analogue of unitaries)

Proof:

To show that  $\underline{\sigma}$  is orthogonal we just need to show that  $|\underline{\sigma}| = |\underline{\sigma}|$

To do so, we note that  $\text{Tr}((\underline{\sigma} \cdot \underline{\sigma})^2) = \text{Tr}((\sum_i \sigma_i \sigma_i)^2)$

$$\begin{aligned}
 &= \text{Tr}(\sum_{ij} s_i s_j \sigma_i \sigma_j) \\
 &= \sum_{ij} s_i s_j \delta_{ij} \\
 &= \sum_i s_i^2 = |\underline{s}|^2
 \end{aligned}$$

Now we note that

$$\begin{aligned}
 |\underline{s}|^2 &= \text{Tr}((\underline{s} \cdot \underline{\sigma})^2) = \text{Tr}((U (\underline{s} \cdot \underline{\sigma}) U^\dagger)^2) = \text{Tr}(U (\underline{s} \cdot \underline{\sigma})^2 U^\dagger) = \sum_{ij} \text{Tr}(\tau_i \tau_j \sigma_i \sigma_j) \\
 &= \sum \tau_i^2 = |\underline{\tau}|^2
 \end{aligned}$$

$$\therefore |\underline{s}| = |\underline{\tau}| \Rightarrow O \text{ is orthogonal}$$

To get a better handle on the nature of this rotation, note that:

Any unitary on a qubit can be written as

$$U = e^{i\Theta \underline{n} \cdot \underline{\sigma}}$$

(if this isn't immediately obvious note that any unitary can be written as  $U = e^{-iHt}$  and  $H$  can always be expanded in the Pauli basis)

$H$  is now a Hamiltonian (sorry!).

which is equivalent to

$$U = \boxed{e^{-i\Theta \underline{n} \cdot \underline{\sigma}}} = \cos(\Theta) \mathcal{I} - i \sin(\Theta) \underline{n} \cdot \underline{\sigma}$$

(if you've not shown this before it's worth proving it to yourself)

What is the effect of evolving  $10^{10}$  under  $e^{-i\theta O_x}$ ?

$$e^{-i\theta \sigma_x} |10\rangle = \cos(\theta) |10\rangle - i \sin(\theta) \sigma_x |10\rangle$$

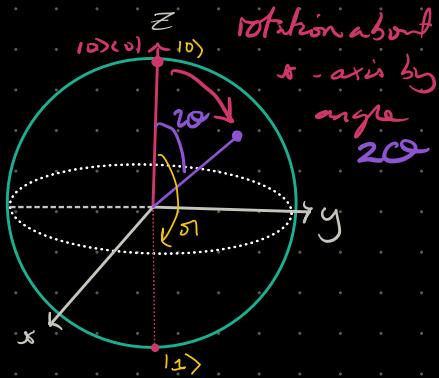
$$= \cos(\theta) |10\rangle - i \sin(\theta) |11\rangle$$

$$\text{Ans. eg. } e^{-\frac{\sigma}{2} \sigma_x} |10\rangle = \cos\left(\frac{\sigma}{2}\right) |10\rangle - i \sin\left(\frac{\sigma}{2}\right) |11\rangle$$

$= -11 + 10$

↑  
unphysical  
global phase

Exercise: Show that  $e^{-i\theta \hat{n} \cdot \vec{v}}$  corresponds to a rotation about  $\vec{n}$  vector by an angle  $2\theta$ .



## ② Two-Qubit Gates

## Examples

- Can just apply multiple copies of a gate onto systems in parallel.

$$\text{eg. } \underbrace{X \otimes X}_{\text{Alternative common notation}} \quad \underbrace{10\rangle \otimes 10\rangle}_{X^{\otimes 2}} = X|0\rangle \otimes X|0\rangle = |2\rangle \otimes |2\rangle$$

- More interesting: CNOT : "Apply NOT to 2nd Qubit iff 1st qubit is in state  $|1\rangle$ "

$$\begin{aligned}
 \text{c.e.} \quad \text{CNOT } |00\rangle &= |00\rangle \\
 \text{CNOT } |01\rangle &= |01\rangle \\
 \text{CNOT } |10\rangle &= |11\rangle \\
 \Rightarrow \text{CNOT } |11\rangle &= |10\rangle
 \end{aligned}$$

Before	After
00	00
01	01
10	11

11 110

Explicit Matrix form:  $CNOT = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$

(Check this really does do)

$CNOT$  can generate entanglement.

eg: consider the sequence of gates

$$\begin{aligned} (CNOT)(H \otimes I)|100\rangle &= CNOT \underbrace{|1+0\rangle}_{\substack{\text{non-classical} \\ \text{but unentangled state}}} \\ &= CNOT \frac{1}{\sqrt{2}}(|100\rangle + |120\rangle) \\ &= \frac{1}{\sqrt{2}}(|100\rangle + |121\rangle) \underbrace{\quad}_{\substack{\text{Bell state}}} \underbrace{\quad}_{\substack{\text{maximally} \\ \text{entangled state}}} \end{aligned}$$

Bell states  $|\phi^\pm\rangle = \frac{1}{\sqrt{2}}(|100\rangle \pm |121\rangle)$  ← with prob  $\frac{1}{2}$  both 00  
 $|\psi^\pm\rangle = \frac{1}{\sqrt{2}}(|101\rangle \pm |120\rangle)$  ← with prob  $\frac{1}{2}$  first qubit 0 & second 1  
probs  $\frac{1}{2}$  first qubit 1 & second qubit 0

How would you prepare the other Bell states using quantum gates?

### ③ Drawing circuits

A tidy way of visualising a series of gates

Best understood by example

